

427(1): Hamilton Equations for m Theory

Consider the relativistic momentum in m theory:

$$p_{||} = \gamma m v_{||} \quad - (1)$$

in frame (r_1, ϕ) , where:

$$r_1 = \frac{r}{m(r_1)^{1/2}} \quad - (2)$$

Here:

$$v_{||} = \frac{v_N}{m(r_1)^{1/2}} \quad - (3)$$

also v_N is the Newtonian velocity.

It follows that:

$$p_{||}^2 c^2 = \gamma^2 m^2 v_{||}^2 c^2 \quad - (4)$$

where

$$\gamma = \left(m(r_1) - \frac{v_{||}^2}{c^2} \right)^{-1/2} \quad - (5)$$

Therefore:

$$p_{||}^2 c^2 = \gamma^2 m^2 c^4 \frac{v_{||}^2}{c^2} \quad - (6)$$

$$= \gamma^2 m^2 c^4 \left(m(r_1) - \frac{1}{\gamma^2} \right)$$

$$= m(r_1) \gamma^2 m^2 c^4 - m^2 c^4$$

$$= \frac{E^2}{m(r_1)} - m^2 c^4$$

It follows that:

$$E^2 = m(r_1) (p_{||}^2 c^2 + m^2 c^4) \quad - (7)$$

as a UFT 424.

The Hamiltonian of m theory is therefore:

2)

$$H = m(r_1) \left(mc^2 - \frac{2MG}{r_1} \right) \quad (8)$$

$$= \left(m(r_1) (p_1^2 c^2 + m^2 c^4) \right)^{1/2} - \frac{2MG}{r_1}$$

Now define the canonically conjugate generalized coordinates:

$$p_r = p_1, \quad q_r = r_1 \quad (9)$$

$$p_\phi = L_1, \quad q_\phi = \phi \quad (10)$$

and

The Hamilton equations are:

$$\dot{p}_r = - \frac{\partial H}{\partial q_r} \quad (11)$$

$$\dot{q}_r = \frac{\partial H}{\partial p_r} \quad (12)$$

$$\dot{p}_\phi = - \frac{\partial H}{\partial q_\phi} \quad (13)$$

$$\dot{q}_\phi = \frac{\partial H}{\partial p_\phi} \quad (14)$$

From eq. (11):

$$\begin{aligned} \dot{p}_1 &= - \frac{2MG}{r_1^2} - \frac{\partial}{\partial r_1} \left(m(r_1) (p_1^2 c^2 + m^2 c^4) \right)^{1/2} \\ &= - \frac{2MG}{r_1^2} - \frac{E}{m(r_1)^{1/2}} \frac{\partial}{\partial r_1} \left(m^{1/2}(r_1) \right) \end{aligned} \quad (15)$$

3) From the Lagrangian analysis of UFT417:

$$P_1 = -\frac{m\dot{m}b}{r_1^2} - mc^2 \frac{dm(r_1)}{dr_1} \quad (16)$$

Now use:

$$E = m(r_1) \gamma mc^2 \quad (17)$$

and compare eqs. (15) and (16):

$$\frac{E}{m(r_1)^{1/2}} \frac{d}{dr_1} (m^{1/2}(r_1)) = \frac{E}{2m(r_1)} \frac{dm(r_1)}{dr_1} \quad (18)$$

$$\text{i.e.} \quad \frac{d}{dr_1} (m^{1/2}(r_1)) = \frac{1}{2m^{1/2}(r_1)} \frac{dm(r_1)}{dr_1} \quad (19)$$

Let

$$f = m(r_1)^{1/2} \quad (20)$$

then

$$\begin{aligned} \frac{df}{dr_1} &= \frac{df}{dm(r_1)} \frac{dm(r_1)}{dr_1} \quad (21) \\ &= \frac{1}{2m^{1/2}(r_1)} \frac{dm(r_1)}{dr_1} \end{aligned}$$

Therefore eq. (21) is true, Q.E.D.

The Hamilton and Euler Lagrange equations produce the same force due to m space:

$$F = -\frac{mc^2 \gamma}{2} \frac{dm(r_1)}{dr_1} = -\frac{E}{m(r_1)^{1/2}} \frac{d}{dr_1} (m^{1/2}(r_1)) \quad (22)$$

4) The Hamilton equation (12) gives:

$$\dot{r}_1 = \frac{\partial H}{\partial p_1} \quad - (23) \quad - (24)$$

also $H = \left(m(r_1) (\dot{p}_1^2 c^2 + m^2 c^4) \right)^{1/2} - \frac{nm\Gamma}{r_1}$

Therefore: $\frac{\partial H}{\partial p_1} = \frac{m(r_1) p_1 c^2}{m(r_1) \gamma m c^2} = \frac{p_1}{m \gamma} \quad - (25)$

also we have used:

$$E = m(r_1) \gamma m c^2 = \left(m(r_1) (\dot{p}_1^2 c^2 + m^2 c^4) \right)^{1/2} \quad - (26)$$

So $p_1 = m \gamma \dot{r}_1 \quad - (27)$

Q.E.D.

To extend to plane polar coordinates use:

$$\underline{p}_1 = \gamma m \underline{\dot{r}}_1 = \gamma m \left(\dot{r}_1 \underline{e}_r + r_1 \dot{\phi} \underline{e}_\phi \right) \quad - (28)$$

So $p_1^2 = p_r^2 + \frac{p_\phi^2}{r^2} \quad - (29)$

$$= p_r^2 + \frac{L^2}{r^2}$$

Eq. (13) gives:

$$\dot{L} = \frac{dL}{dt} = 0 \quad - (30)$$

So L is a constant of motion.

Eq. (14) gives:

$$\dot{\phi} = \frac{\partial H}{\partial L_1} \quad - (31)$$

where

$$H = \left(m(r_1) \left(\left(p_r^2 + \frac{L^2}{r^2} \right) c^2 + m^2 c^4 \right) \right)^{1/2} - \frac{mMG}{r_1}$$

It follows that:

$$\dot{\phi} = \frac{\partial H}{\partial L} = \frac{m(r_1) L_1 c^2}{r_1^2 m(r_1) \gamma m c^2} \quad - (32)$$

$$L_1 = \gamma m r_1^2 \dot{\phi} \quad - (34)$$

Q.E.D.

The n theory in Hamiltona and Euler-Lagrange formalisms is rigorously self consistent, both formalisms give the same result.

The Hamiltona Jacobi equations we found from eq. (32) using:

$$p_r = \frac{\partial S}{\partial r}, \quad p_\phi = \frac{\partial S}{\partial \phi} \quad - (35)$$

where S is the action:

$$S = S_r + S_\phi \quad - (36)$$

Therefore the Hamiltona Jacobi equations of n theory

b) are:

$$E_1 = \left(m(r_1) \left(\left(\frac{dS_r}{dr} \right)^2 + \frac{L^2}{r^2} \right) c^2 + m^2 c^4 \right)^{1/2} - \frac{n h k}{r_1} \quad - (37)$$

and

$$L_1 = \frac{dS_\phi}{d\phi} = \text{constant} \quad - (38)$$

These equations can be integrated by computer algebra to give S_r and S_ϕ . This is a route to quantization because h is the quantum of action.

The spacetime can be found for eq. (22)
