

4/16(2): The Self Consistent Coordinate System
 Consider the plane polar coordinate system (r, ϕ) ,

line:
$$r_1 = \frac{r}{m(r)^{1/2}} \quad (1)$$

The Lagrangian is then:
$$L = -mc^2 \left(m(r) - \frac{1}{c^2} (\dot{r}_1^2 + r_1^2 \dot{\phi}^2) \right)^{1/2} + \frac{m\hbar\omega}{r_1} \quad (2)$$

where
$$\dot{r}_1 = \frac{\dot{r}}{m(r)^{1/2}} \quad (3)$$

The Lagrangian has the fundamental property:

$$\frac{\partial L}{\partial t} = 0 \quad (4)$$

(Morse and Feshbach). The linear momentum is:
$$\underline{p}_1 = \frac{\partial L}{\partial \dot{r}_1} = \gamma m \dot{r}_1 = \frac{\gamma m \dot{r}}{m(r)^{1/2}} \quad (5)$$

This result is also obtained from kinematics w.r.t.:

$$\underline{r} = \frac{r}{m(r)^{1/2}} \underline{e}_r \quad (6)$$

and as derived in 4/17/15 from fundamental geometry.

The angular momentum for eq. (2) is:

$$L = \frac{\partial L}{\partial \dot{\phi}} = \gamma m r_1^2 \dot{\phi} = \frac{\gamma m r^2 \dot{\phi}}{m(r)} \quad (7)$$

From eqs. (5) and (6):

$$\underline{L} = \underline{r} \times \underline{p} = \frac{\gamma m r^2 \dot{\phi}}{m(r)} \underline{k} \quad (8)$$

s. of them is self consistent, Q.E.D.

From Lagrangian dynamics, the Leibniz equation in
 - space is:

$$\underline{\dot{p}}_1 = \frac{\partial \mathcal{L}}{\partial \underline{r}_1} \quad - (9)$$

i.e.

$$\frac{d}{dt} \left(\frac{\gamma m \underline{\dot{r}}}{m(r)^{1/2}} \right) = - \frac{n m b}{r_1^2} \quad - (10)$$

i.e.

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \underline{\dot{r}}_1} \right) = \frac{\partial \mathcal{L}}{\partial \underline{r}_1} \quad - (11)$$

which is the Euler Lagrange equation in n space.
 Note that the Lagrangian (2) can be written as:

$$\mathcal{L} = -m c^2 \left(m(r) - \frac{1}{c^2 m(r)} (\dot{r}^2 + r^2 \dot{\phi}^2) \right)^{1/2} + \frac{m^{1/2}(r) n m b}{r} \quad - (12)$$

so the potential energy is changed to:

$$U(r) = - \frac{n m b}{r} = - m^{1/2}(r) \frac{n m b}{r} \quad - (13)$$

In a curved spacetime such as n spacetime the
 distance between n and m is $r / m(r)$. For
 small curvature:

$$r_1 \rightarrow r \quad - (14)$$

From Euler Lagrange equation:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) = \frac{\partial \mathcal{L}}{\partial \phi} \quad - (15)$$

it follows that \mathcal{L} is conserved:

$$\frac{d\mathcal{L}}{dt} = 0 \quad - (16)$$

is UFT415.

The only difference between the Hamiltonian and Lagrangian used in UFT415 is that the polar coordinate system is changed to (r, ϕ) . In the coordinate system there is no need for an $m^{1/2}(r)$ factor in eq. (57) of UFT415.

In Hamiltonian dynamics the canonical equations

$$\dot{r} = \frac{\partial H}{\partial p} \quad (17)$$

$$\dot{p} = - \frac{\partial H}{\partial r} \quad (18)$$

where H is the Hamiltonian. In the (r, ϕ) coordinate system these are changed to:

$$\dot{r}_1 = \frac{\partial H}{\partial p_1} \quad (19)$$

$$\dot{p}_1 = - \frac{\partial H}{\partial r_1} \quad (20)$$

In Hamiltonian dynamics:

$$H = H(r, p, t) \quad (21)$$

In Lagrangian dynamics:

$$L = L(r, \dot{r}, t) \quad (22)$$

Mercer and Thorne).

$$\text{Using } H = m(r) \gamma mc^2 - \frac{mMG}{r_1} \quad (23)$$

$$= m(r) \left(m(r) - \frac{v_1^2}{c^2} \right)^{-1/2} mc^2 - \frac{mMG}{r_1}$$

it follows that

$$\dot{r}_1 = \frac{\partial H}{\partial p_1} = \frac{p_1}{m^2 c^2} \gamma m(r) m c^2 = m(r) \gamma \frac{p_1}{m} \quad (24)$$

also

$$p_1 = \frac{v}{m(r)^{1/2}} \quad (25)$$

so

$$v = \frac{\dot{r}_1}{m(r)^{1/2}} \quad (26)$$

Q.E.D.

Eq. (26) is the kinematic result used in 4FT415,
the constant of motion:

$$H = m(r) \gamma m c^2 - \frac{2mG}{r_1} \quad (27)$$

self consistently defined, Q.E.D.

The Hamiltonian H has the property:

$$\frac{dH}{dt} = 0 \quad (28)$$

Similarly the Lagrangian (2) has the property:

$$\frac{dL}{dt} = 0 \quad (29)$$
