

4/10(1): The Correct Theory of de Sitter Precession

The standard model of de Sitter precession is the rotation of an infinitesimal line element of \mathcal{Q} type:

$$ds^2 = n(r, t) c^2 dt^2 - \frac{dr^2}{n(r, t)} - r^2 d\phi^2 \quad (1)$$

in polar coordinates, where $n(r, t)$ is a function of r and t . The ECE2 metric is characterized by:

$$n(r, t) = 1 \quad (2)$$

and the so-called Schwarzschild metric by:

$$n(r, t) = 1 - \frac{2mG}{r} = 1 - \frac{r_0}{r} \quad (3)$$

The rotation is defined for all $n(r, t)$ by:

$$\phi' = \phi + \omega t \quad (4)$$

and

$$v_\phi = \omega r \quad (5)$$

The rotation (4) is the de Sitter rotation, and was originally applied in 1916 to the Schwarzschild metric, defined by eq. (3). It is now known that this procedure is entirely incorrect because it is based on an incorrect geometry without basis (UFT88, UFT99, UFT313). The Wikipedia article on this subject is entirely obscure, so is the entire correct mathematics of de Sitter precession are given in order to complete the result with the ECE2 theory.

From eqs. (1) to (5), following the method of UFT409:

$$ds^2 = (n(r, t) c^2 - v_\phi^2) dt^2 - \frac{dr^2}{n(r, t)} - r^2 d\phi^2 - 2\omega r^2 d\phi dt \quad (6)$$

in which

$$d\phi = \omega dt \quad (7)$$

so it follows that:

$$1) ds^2 = (m(r,t)c^2 - 3v_\phi^2) dt^2 - \left(\frac{dr^2}{m(r,t)} + r^2 d\phi^2 \right) \quad (8)$$

Define:
$$v_1^2 = \frac{1}{m(r,t)} \left(\frac{dr}{dt} \right)^2 + r^2 \left(\frac{d\phi}{dt} \right)^2 \quad (9)$$

to find that
$$ds^2 = (m(r,t)c^2 - 3v_\phi^2 - v_1^2) dt^2 \quad (10)$$

$$= c^2 d\tau^2$$

$$= \left(m(r,t) - \frac{(3v_\phi^2 + v_1^2)}{c^2} \right) c^2 dt^2$$

It follows that:
$$d\tau = \left(m(r,t) - \frac{v^2}{c^2} \right)^{1/2} dt \quad (11)$$

Now define:
$$\gamma_1 = \frac{dt}{d\tau} = \left(m(r,t) - \frac{v^2}{c^2} \right)^{-1/2} \quad (12)$$

where
$$v^2 = 3v_\phi^2 + v_1^2 \quad (13)$$

Note that eq. (12) reduces to the Lorentz factor:

$$\gamma = \frac{dt}{d\tau} = \left(1 - \frac{v^2}{c^2} \right)^{-1/2} \quad (14)$$

where:
$$m(r,t) = 1 \quad (15)$$

and
$$v_\phi = 0 \quad (16)$$

Time dilatation is defined by:

$$dt = \gamma_1 d\tau \quad (17)$$

Note that the line element is defined by:

$$ds^2 = \left(n(r,t) - \frac{v}{c} \right)^2 c^2 dt^2 - (18)$$

$$= c^2 dt_1^2$$

in the observer frame. By definition:

$$dt_1 = dt - (19)$$

Note that

$$dt_1 < dt - (20)$$

and that a precession can be defined by:

$$\Delta \phi = \omega_0 (dt - dt_1) - (21)$$

For a rotation of 2π :

$$\omega_0 dt_1 := 2\pi - (22)$$

So the de Sitter precession is:

$$\Delta \phi_g = 2\pi (\gamma - 1) - (23)$$

This is the curved calculation of the standard

model's de Sitter precession.

It reduces to the ECE precession when:

$$n(r,t) = 1. - (24)$$

When there is no de Sitter rotation, eq. (23) reduces

$$\Delta \phi = 2\pi (\gamma - 1) - (25)$$

to where γ is the Lorentz factor:

$$\gamma = \left(1 - \frac{v^2}{c^2} \right)^{-1/2} - (26)$$

The precession (25) occurs in the unrotated $E(E)$ element:

$$ds^2 = c^2 d\tau^2 = (c^2 - v_w^2) dt^2 \quad (27)$$

ω is defined by:

$$\Delta \phi = \omega_0 (dt - d\tau) \quad (28)$$

$$= 2\pi (\gamma - 1)$$

for a rotation of 2π .

Time dilation or dilatation, is defined as

$$dt = \gamma d\tau \quad (29)$$

so:

$$dt > d\tau \quad (30)$$

and the time interval in the observer frame dt is larger than in the moving frame, the proper time interval $d\tau$. This is the traditional interpretation of special relativity, but eq. (19) shows that $d\tau$ can be thought of as interval dt_1 in the observer frame, so dt and dt_1 are the same, observer, frame.

There is therefore only one source of precession, difference between dt and dt_1 .

The only correct source of precession is the $E(E)$ theory, because this does not rely on the Einstein old equation. Furthermore it is not possible to isolate the various contributions to precession, as explained in UFT46.