

580 (1): Evaluation of  $\underline{\Phi}$  and Spin Connection Vectors  
 From Note 379 it was found that the antisymmetry condition  
 reduces the following four equations in the limit:

$$\underline{\Phi} = \underline{MG} \quad - (1)$$

The first two are force equations:

$$\ddot{x} = -\frac{mG}{(x^2+y^2)^{3/2}} \left( \frac{x}{x^2+y^2} - \omega_x \right) \quad - (2)$$

$$\ddot{y} = -\frac{mG}{(x^2+y^2)^{3/2}} \left( \frac{y}{x^2+y^2} - \omega_y \right) \quad - (3)$$

where the spin connection vector is:

$$\underline{\omega} = \omega_x \underline{i} + \omega_y \underline{j} \quad - (4)$$

By antisymmetry:

$$\left( \frac{d}{dt} + \omega_0 \right) a_x = \frac{mG}{(x^2+y^2)^{3/2}} \left( \frac{x}{x^2+y^2} - \omega_x \right) \quad - (5)$$

$$\left( \frac{d}{dt} + \omega_0 \right) a_y = \frac{mG}{(x^2+y^2)^{3/2}} \left( \frac{y}{x^2+y^2} - \omega_y \right) \quad - (6)$$

There are five unknowns,  $\omega_0$ ,  $\omega_x$ ,  $\omega_y$ ,  $a_x$  and  $a_y$ ,  
 and to determine these uniquely requires more information.  
 The gauge-invariant field equation is:

$$\underline{\nabla} \cdot \underline{\Omega} = 0 \quad - (7)$$

Assumed absence of gauge-invariant monopole.

\* The gravitomagnetic field is found from:

$$b_{12} = (d_1 + \omega_1) Q_2 - (d_2 + \omega_2) Q_1 \quad - (8)$$

$$b_{13} = (d_1 + \omega_1) Q_3 - (d_3 + \omega_3) Q_1 \quad - (9)$$

$$b_{23} = (d_2 + \omega_2) Q_3 - (d_3 + \omega_3) Q_2 \quad - (10)$$

where the field tensor is:

$$b_{\mu\nu} = \begin{bmatrix} 0 & -g_1/c & -g_2/c & -g_3/c \\ g_1/c & 0 & \Omega_3 & -\Omega_2 \\ g_2/c & -\Omega_3 & 0 & \Omega_1 \\ g_3/c & \Omega_2 & -\Omega_1 & 0 \end{bmatrix} \quad - (11)$$

Therefore

$$\underline{\Omega} = \underline{\nabla} \times \underline{Q} - \underline{\omega} \times \underline{Q} \quad - (12)$$

From eqs. (7) and (12)

$$\underline{\nabla} \cdot \underline{\omega} \times \underline{Q} = \underline{Q} \cdot \underline{\nabla} \times \underline{\omega} - \underline{\omega} \cdot \underline{\nabla} \times \underline{Q} = 0 \quad - (13)$$

The antisymmetry law for eqs. (8) to (10) is:

$$(d_1 + \omega_1) Q_2 = - (d_2 + \omega_2) Q_1 \quad - (14)$$

$$(d_1 + \omega_1) Q_3 = - (d_3 + \omega_3) Q_1 \quad - (15)$$

$$(d_2 + \omega_2) Q_3 = - (d_3 + \omega_3) Q_2 \quad (16)$$

For planar orbital motion:

$$(d_1 + \omega_1) Q_2 = - (d_2 + \omega_2) Q_1 \quad - (17)$$

3) so 
$$\left(\frac{\partial}{\partial x} - \omega_x\right) Q_y = - \left(\frac{\partial}{\partial y} - \omega_y\right) Q_x - (18)$$

for planar motion is general.

Eqs. (2), (3), (5), (6) and (18) are five equations in five unknowns,  $\omega_x, \omega_y, Q_x$  and  $Q_y$ . This set of equations can be solved numerically.

In addition, the orbit from eqns. (2) and (3) must be the same as the orbit from the Hamiltonian:

$$H = \gamma mc^2 + U - (19)$$

and the Lagrangian:

$$L = -\frac{mc^2}{\gamma} - U - (20)$$

The solution for  $\underline{\omega}$  and  $\underline{Q}$  must obey the

equation: 
$$\underline{\nabla} \cdot \underline{\omega} \times \underline{Q} = \underline{0} - (21)$$

Eq. (21) is obeyed by planar motion in vcd:

$$\underline{\omega} = \omega_x \underline{i} + \omega_y \underline{j} - (22)$$

$$\underline{Q} = Q_x \underline{i} + Q_y \underline{j} - (23)$$

if 
$$\frac{\partial}{\partial t} (\omega_x Q_y - \omega_y Q_x) = 0 - (24)$$

This is true automatically if  $\underline{\omega}$  and  $\underline{Q}$  are

1) Defining a plane  $XY$  and have no  $Z$  dependence.

In general the ECE2 gravitational field equations imply that  $\underline{\Omega}$  exists if  $g$  exists. The components of the gravitomagnetic field are in general:

$$\Omega_1 = (\partial_2 + \omega_2) a_3 - (\partial_3 + \omega_3) a_2 \quad (18)$$

$$\Omega_2 = (\partial_3 + \omega_3) a_1 - (\partial_1 + \omega_1) a_3 \quad (19)$$

$$\Omega_3 = (\partial_1 + \omega_1) a_2 - (\partial_2 + \omega_2) a_1 \quad (20)$$

For a planar orbit:

$$\Omega_3 = (\partial_1 + \omega_1) a_2 - (\partial_2 + \omega_2) a_1 \quad (21)$$

i.e.

$$\Omega_z = \left( \frac{d}{dx} - \omega_x \right) a_y - \left( \frac{d}{dy} - \omega_y \right) a_x \quad (22)$$

so the gravitomagnetic field is perpendicular to the plane of the orbit. From eqs. (18) and (22):

$$\begin{aligned} \Omega_z &= 2 \left( \frac{d}{dx} - \omega_x \right) a_y \\ &\quad - 2 \left( \frac{d}{dy} - \omega_y \right) a_x \end{aligned} \quad (23)$$

Newtonian Limit

In this limit:

$$\omega_0 = 0 \quad (24)$$

$$\omega_x = \omega_y = 0 \quad (25)$$

so:

$$\ddot{x} = -\frac{\partial Q_x}{\partial t} = \frac{mGx}{(x^2+y^2)^{3/2}} \quad - (26)$$

$$\text{and } \ddot{y} = -\frac{\partial Q_y}{\partial t} = \frac{mGy}{(x^2+y^2)^{3/2}} \quad - (27)$$

Eq. (18) reduces to:

$$\frac{\partial Q_y}{\partial x} = -\frac{\partial Q_x}{\partial y} \quad - (28)$$

$$\text{and } \Omega_z = 2\frac{\partial Q_y}{\partial x} = -2\frac{\partial Q_x}{\partial y} \quad - (29)$$

because in Newtonian gravitation there is no gravito-magnetic field. The orbit from eqs. (26) and (27) is an ellipse. In the presence of a spin correction the orbit is a forward or retrograde precessing ellipse.

Zero Gravitation Limit

In this case:

$$\ddot{x} = \ddot{y} = 0 \quad - (30)$$

$$\text{so: } \omega_x = \frac{x}{x^2+y^2} \quad - (31)$$

$$\omega_y = \frac{y}{x^2+y^2} \quad - (32)$$

Therefore eqs. (5) and (6) reduce to:

$$\left(\frac{d}{dt} + \omega_0\right) Q_x = \left(\frac{d}{dt} + \omega_0\right) Q_y = 0 \quad - (33)$$

6) Eq. (18) reduces to:

$$\left(\frac{d}{dt} - \frac{x}{x^2 + y^2}\right) Q_y = - \left(\frac{d}{dt} - \frac{y}{x^2 + y^2}\right) Q_x \quad (34)$$

From eq. (33):  $Q_x = Q_y \quad (35)$

so there are two equations in two unknowns:

$$\left(\frac{d}{dt} + \omega_0\right) Q_x = 0 \quad (36)$$

and  $\left(\frac{d}{dt} - \frac{x}{x^2 + y^2}\right) Q_x = - \left(\frac{d}{dt} - \frac{y}{x^2 + y^2}\right) Q_x \quad (37)$

So any reports of zero gravitation can be analyzed in this way. Counter gravitation means changing the signs of  $x$  and  $y$  in eqs. (2) and (3).

The above analysis depends on assumption (1), i.e. that the scalar potential is essentially Newtonian, and that orbital precessions are small. More generally:

$$\underline{g} = - \underline{\nabla} \underline{\Phi} + \underline{\Phi} \underline{\omega} = - \frac{\partial \underline{Q}}{\partial t} - \underline{\omega}_0 \underline{Q} \quad (38)$$

$$\underline{\nabla} \cdot \underline{g} = \underline{\kappa} \cdot \underline{g} = 4\pi G \rho_m \quad (39)$$

$$\underline{\Omega} = \underline{\nabla} \times \underline{Q} - \underline{\omega} \times \underline{Q} \quad (40)$$

$$\underline{\nabla} \cdot \underline{\Omega} = 0 \quad (41)$$