

4(4): Resonance Condition for Gravitational Potentials

The ECE wave equation is:

$$\nabla^2 \Phi = R \Phi \quad - (1)$$

Let Φ is the gravitational potential and R the scalar curvature. The solution of Eq. (1) is:

$$\Phi = \Phi_0 \exp(i(\omega t - \underline{\kappa} \cdot \underline{r})) \quad - (2)$$

also:

$$\frac{\omega^2}{c^2} - \kappa^2 = R \quad - (3)$$

For geostatics it is assumed that:

$$\Phi = \Phi_0 \exp(-i \underline{\kappa}_0 \cdot \underline{r}) \quad - (4)$$

and

$$(\nabla^2 + R) \Phi = 0 \quad - (5)$$

also

$$R = \kappa_0^2 \quad - (6)$$

A small driving force of electromagnetic origin is applied to the gravitational system consisting of a mass m attracted by a mass M . Therefore Eq. (5) becomes the Helmholtz equation:

$$(\nabla^2 + R) \Phi = A \cos \underline{\kappa} \cdot \underline{r} \quad - (7)$$

In one dimension:

$$\left(\frac{\partial^2}{\partial z^2} + R \right) \Phi = A \cos \kappa_z z \quad - (8)$$

The electromagnetic driving force is applied along the z axis.

The solution of Eq. (8) is:

$$\underline{\Phi} = \frac{A \cos k_z z}{k_0^2 - k_z^2} \quad - (9)$$

and counter gravitation occurs under the condition:

$$k_0 = k_z \quad - (10)$$

when $\underline{\Phi} \rightarrow \infty \quad - (11)$

In ECE2:

$$\underline{g} = -\underline{\nabla} \underline{\Phi} - \frac{d\underline{Q}}{dt} \quad - (12)$$

and

$$\underline{\nabla} \cdot \underline{g} = 4\pi G \rho_m \quad - (13)$$

where \underline{Q} is the gravitational vector potential. This does not exist in the standard model of physics. Here G is Newton's constant and ρ_m is the source mass density.

So

$$-\nabla^2 \underline{\Phi} - \underline{\nabla} \cdot \frac{d\underline{Q}}{dt} = 4\pi G \rho_m \quad - (14)$$

If it is assumed that:

$$\underline{Q} = \underline{0} \quad - (15)$$

then:

$$\nabla^2 \underline{\Phi} = -4\pi G \rho_m \quad - (16)$$

From eqs. (7) and (16):

$$-4\pi G \rho_m + k_0^2 \underline{\Phi} = A \cos k \cdot \underline{r} \quad - (17)$$

In one dimension:

$$-4\pi G \rho_m + \kappa_0^2 \Phi = A \cos \kappa_z z \quad - (18)$$

m eqs. (9) and (18):

$$-\frac{4\pi}{\lambda} G \rho_m + \kappa_0^2 A \cos \kappa_z z = A \cos \kappa_z z \quad - (19)$$

$$\therefore \kappa_0^2 A \cos \kappa_z z = \frac{\kappa_0^2 - \kappa_z^2}{\kappa_0^2 - \kappa_z^2} (A \cos \kappa_z z + \frac{4\pi}{\lambda} G \rho_m) \quad - (20)$$

From eqs. (10) and (20) the resonance condition is

$$\cos \kappa_z z = 0 \quad - (21)$$

i.e.

$$\kappa_z z = \frac{\pi}{2} \quad - (22)$$

Finally assume that:

$$\kappa_z = \frac{\omega}{c} \quad - (23)$$

then:

$$\omega z = \frac{\pi}{2} c \quad - (24)$$

If z is one metre then: - (25)

$$\omega = 4.709 \times 10^8 \text{ rad/sec}$$

This result seems to be too simple, because mass held one metre above the surface of the earth is not elevated by irradiating it by e/ν radiation.

The origin of the problem is that Φ has been

assumed to be a wave in eq. (1). Hence it has been assumed that there exist gravitational waves. In order for such waves to exist the gravitational vector potential \underline{Q} has to be non-zero, and the gravitational field must exist. Additionally, in analogy with the Maxwell displacement current, gravitational waves need the existence of the "Maxwell displacement current" of $E(E_2)$, i.e. $\frac{dQ}{dt}$.

So the two relevant equations are:

$$(\nabla^2 + R)\underline{\Phi} = A \cos \underline{\kappa} \cdot \underline{r} \quad - (26)$$

and

$$\nabla^2 \underline{\Phi} + \underline{\nabla} \cdot \frac{d\underline{Q}}{dt} = -4\pi G \rho_m \quad - (27)$$

Hence:

$$R \underline{\Phi} - \underline{\nabla} \cdot \frac{d\underline{Q}}{dt} - 4\pi G \rho_m = A \cos \underline{\kappa} \cdot \underline{r} \quad - (28)$$

In one dimension:

$$\frac{\kappa_0^2 A \kappa_z Z}{\kappa_0^2 - \kappa_z^2} - 4\pi G \rho_m - \frac{dQ}{dt} = A \cos \kappa_z Z \quad - (29)$$

$$\kappa_0^2 A \kappa_z Z = (\kappa_0^2 - \kappa_z^2) \left(A \cos \kappa_z Z + \frac{dQ}{dt} + 4\pi G \rho_m \right)$$

However, it is clearer to use:

$$\nabla^2 \underline{\Phi} = -\kappa_z^2 \underline{\Phi} \quad - (30)$$

in eq. (27), so:

$$-\kappa_z^2 \underline{\Phi} + \underline{\nabla} \cdot \frac{\underline{\partial Q}}{\underline{\partial t}} = -\frac{4\pi \underline{G} \rho}{m} \quad (31)$$

where

$$\underline{\Phi} = \frac{A \cos \kappa_z z}{\kappa_0^2 - \kappa_z^2} \quad (32)$$

$$\text{so } \frac{A \kappa_z^2 \cos \kappa_z z}{\kappa_0^2 - \kappa_z^2} = \underline{\nabla} \cdot \frac{\underline{\partial Q}}{\underline{\partial t}} + \frac{4\pi \underline{G} \rho}{m} \quad (33)$$

so at resonance:

$$\underline{\nabla} \cdot \frac{\underline{\partial Q}}{\underline{\partial t}} \rightarrow \infty \quad (34)$$

when:

$$\kappa_0 = \kappa_z \quad (35)$$

and

$$\underline{\Phi} \rightarrow \infty \quad (36)$$

In order for this method to work the electromagnetic frequency ω_z must be tuned to a characteristic frequency $\omega_0 = \kappa_0 c$.

If it is assumed that \underline{g} is not a wave,

$$\text{then } \underline{g} = -\underline{\nabla} \underline{\Phi} \quad (37)$$

More generally,

$$\underline{g} = -\underline{\nabla} \underline{\Phi} - 2\underline{\Phi} \underline{\omega} + 2c\omega_0 \underline{Q} - \frac{\underline{\partial Q}}{\underline{\partial t}} \quad (38)$$