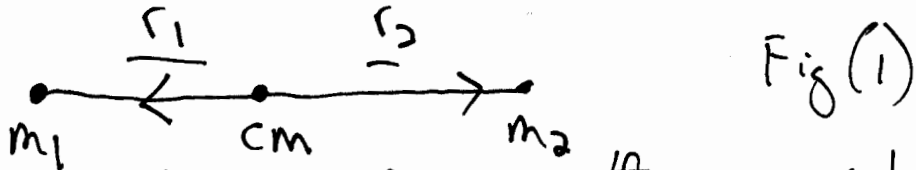


375(8): Relativistic Lagrangian for the Binary Pulsar

If the masses of the two stars are m_1 and m_2 , then we refer to Fig (1):



where cm denotes the centre of mass, the non-relativistic Lagrangian is:

$$L = \frac{1}{2} m_1 \dot{\underline{r}}_1 \cdot \dot{\underline{r}}_1 + \frac{1}{2} m_2 \dot{\underline{r}}_2 \cdot \dot{\underline{r}}_2 + \frac{m_1 m_2 G}{r} \quad (1)$$

where

$$r = |\underline{r}_1 - \underline{r}_2| = \left((\underline{r}_1 - \underline{r}_2) \cdot (\underline{r}_1 - \underline{r}_2) \right)^{1/2} \quad (2)$$

In Cartesian coordinates:

$$\underline{r}_1 = x_1 \underline{i} + y_1 \underline{j} \quad (3)$$

$$\underline{r}_2 = x_2 \underline{i} + y_2 \underline{j} \quad (4)$$

So:

$$L = \frac{1}{2} m_1 (\dot{x}_1^2 + \dot{y}_1^2) + \frac{1}{2} m_2 (\dot{x}_2^2 + \dot{y}_2^2) + \frac{m_1 m_2 G}{r} \quad (5)$$

where

$$r = \left((x_1 - x_2)^2 + (y_1 - y_2)^2 \right)^{1/2} \quad (6)$$

The proper Lagrange variables are x_1, y_1, x_2 and y_2 .

2) so: $\frac{\partial \mathcal{L}}{\partial x_1} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}_1}$ - (7)

$\frac{\partial \mathcal{L}}{\partial y_1} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{y}_1}$ - (8)

$\frac{\partial \mathcal{L}}{\partial x_2} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}_2}$ - (9)

$\frac{\partial \mathcal{L}}{\partial y_2} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{y}_2}$ - (10)

These four simultaneous equations can be integrated numerically to give the orbits of m_1 and m_2 around the centre of mass. These are ellipses with the centre of mass as one focus of the ellipse. This calculation checks the Wegatia code.

The earth-sun system is described by:

$m_2(\text{sun}) \gg m_1(\text{earth})$ - (11)

and $r_1 \gg r_2$ - (12)

so the centre of mass is essentially the position of the sun. Then eq. (5) reduces to:

$\mathcal{L}(\text{earth/sun}) = \frac{1}{2} m_1 (\dot{x}_1^2 + \dot{y}_1^2) + \frac{m_1 m_2 G}{(x_1^2 + y_1^2)^{1/2}}$ - (13)

with the Euler-Lagrange equations (7) and (8). Eqs. (7), (8) and (13) give an ellipse with m_2 as one focus.

3) This ellipse is, in plane polar coordinates:

$$r = \frac{d}{1 + e \cos \phi} \quad (14)$$

where

$$d = a(1 - e^2) \quad (15)$$

Here a is the semi major axis, e the eccentricity and d the half right distance. If b is the semi minor axis then

$$d = b(1 - e^2)^{1/2} \quad (16)$$

The ellipse from eqs. (7), (8) and (13) should obey:

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1 \quad (17)$$

where:

$$x = c + r \cos \phi \quad (18)$$

$$y = r \sin \phi$$

and

$$e = \left(1 - \frac{b^2}{a^2}\right)^{1/2} = \frac{c}{a} \quad (19)$$

The perihelion of the ellipse is defined by:

$$\phi = 0, \quad \cos \phi = 1 \quad (20)$$

so

$$r_{\min} = \frac{d}{1 + e} \quad (21)$$

The orbital velocity at the perihelion is:

$$v^2 = m_2 G \left(\frac{2}{r_{\min}} - \frac{1}{a} \right) \quad (22)$$

$$= \left(\dot{x}_1^2 + \dot{y}_1^2 \right)^{1/2}$$

4) Therefore at the perihelion:

$$v = \frac{m_2 b}{d} (2(1+\epsilon) + 1 - \epsilon^2) \quad (23)$$

So at the perihelion the negative program should produce:

$$\dot{X}_1^2 + \dot{Y}_1^2 = \frac{m_2 b}{d} (3 + 2\epsilon - \epsilon^2) \quad (24)$$

Relativistic Lagrangian

The relativistic Lagrangian of the binary pair is:

$$\mathcal{L} = -m_1 c^2 \left(1 - \frac{\dot{X}_1^2 + \dot{Y}_1^2}{c^2} \right)^{1/2} - m_2 c^2 \left(1 - \frac{\dot{X}_2^2 + \dot{Y}_2^2}{c^2} \right)^{1/2} + \frac{m_1 m_2 b}{(\dot{X}_1^2 + \dot{Y}_1^2)^{1/2}} \quad (25)$$

and the Euler Lagrange equations (7) to (10) are solved with this Lagrangian.

In the limit:

$$m_2 \gg m_1; \quad r_1 \gg r_2 \quad (26)$$

Eq. (25) reduces to:

$$\mathcal{L} = -m_1 c^2 \left(1 - \frac{\dot{X}_1^2 + \dot{Y}_1^2}{c^2} \right)^{1/2} + \frac{m_1 m_2 b}{(\dot{X}_1^2 + \dot{Y}_1^2)^{1/2}} \quad (27)$$

5) Eqs. (7), (8) and (27) produce a precessing ellipse, defined by $x_1(t)$, $y_1(t)$, $\dot{x}_1(t)$ and $\dot{y}_1(t)$. Similarly eqs. (7) to (10) and (25) produce a combination of precessing ellipses, defined by $x_1(t)$, $y_1(t)$, $x_2(t)$ and $y_2(t)$ and $\dot{x}_1(t)$, $\dot{y}_1(t)$, $\dot{x}_2(t)$ and $\dot{y}_2(t)$. The precessions are maximised by large masses m_1 and m_2 , as it is a binary pulsar.

In a static ellipse the function (14) is the same after:

$$\phi \rightarrow \phi + 2\pi \quad (28)$$

because:

$$\cos \phi = \cos(\phi + 2\pi) \quad (29)$$

but it is a precessing ellipse:

$$\phi \rightarrow \phi + \Delta\phi \quad (30)$$

after a 2π revolution, where $\Delta\phi$ is the precession per orbit. Hence it is a precessing ellipse:

$$r = \frac{d}{1 + \epsilon \cos \phi} \rightarrow \frac{d}{1 + \epsilon \cos(\phi + \Delta\phi)} \quad (31)$$

after one orbit (a 2π revolution of ϕ). So at the perihelion:

$$r_{\min} = \frac{d}{1 + \epsilon} \rightarrow \frac{d}{1 + \epsilon \cos(2\pi + \Delta\phi)} \quad (32)$$

2) In eq. (32):

$$\begin{aligned}\cos(2\pi + \Delta\phi) &= \cos 2\pi \cos \Delta\phi - \sin(2\pi) \sin \Delta\phi \\ &= \cos(\Delta\phi) \quad - (33)\end{aligned}$$

So after one orbit:

$$r_{\text{min}} = \frac{a}{1+\epsilon} \rightarrow \frac{a}{1+\epsilon \cos \Delta\phi} \quad - (34)$$

Suggested Computational Work

- 1) Compute the elliptical orbit for eqs. (7), (8) and (13).
 - 2) Compute & precessing elliptical orbit for eqs. (7), (8) and (27).
 - 3) Measure $\Delta\phi$ graphically and show that the precession increases for a very large $m_2 \gg m_1$.
 - 4) Compute the elliptical orbit for eqs. (7) - (10) and Eq. (5). Demonstrate graphically the effect of increasing m_1 in relation to m_2 .
 - 5) Compute & precessing elliptical orbit for eqs. (7) to (10) and eq. (25). Show that the precession increases as m_1 and m_2 get larger.
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