

233(4) : True Precessing Ellipse from \mathcal{L} (rotated Metric).

The true precessing ellipse is:

$$r = \frac{d}{1 + e \cos(x\theta)} \quad - (1)$$

also:

$$d = (1+e)r_{\min} = (1-e)r_{\max} \quad - (2)$$

Here d is the half right ascension, r_{\min} and r_{\max} the distance of closest approach and furthest separation, e the ellipticity, x the precession factor and (r, θ) the plane polar coordinate system. From eq. (1):

$$\left(\frac{dr}{d\theta}\right)^2 = \left(\frac{ex}{d}\right)^2 r^4 \sin^2(x\theta) \quad - (3)$$

$$= \left(\frac{ex}{d}\right)^2 r^4 \left(1 - \frac{1}{e^2} \left(\frac{d}{r} - 1\right)^2\right)$$

$$= x^2 \left(\frac{r}{d}\right)^2 \left(e^2 r^2 - (d-r)^2\right)$$

$$= x^2 r^2 \left(\frac{1-e}{1+e}\right) \left(\frac{(r_{\max}-r)(r-r_{\min})}{r_{\min}^2}\right)$$

$$= x^2 r^2 \left(\frac{1+e}{1-e}\right) \left(\frac{(r_{\max}-r)(r-r_{\min})}{r_{\max}^2}\right)$$

In the Newtonian limit:

$$x = 1 \quad - (4)$$

2) As in UFT 195, the (rotten) metric is:

$$ds^2 = c^2 d\tau^2 = AC^{1/2} c^2 dt^2 - BC^{1/2} dr^2 - C(r) d\theta^2 \quad - (5)$$

where $C(r) = (|r-r_0|^n + d_1^n)^{2/n} \quad - (6)$

This gives:

$$E = mc^2 AC^{1/2}(r) \left(\frac{dt}{d\tau} \right) \quad - (7)$$

$$L = m C(r) \frac{d\theta}{d\tau} \quad - (8)$$

$$\frac{dt}{d\tau} = \left(AC^{1/2}(r) - \frac{v^2}{c^2} \right)^{-1/2} \quad - (9)$$

and

$$\left(\frac{dr}{d\theta} \right)^2 = \frac{m C(r)}{BL^2} \left(\frac{E^2}{Amc^2} - C^{1/2}(r) \left(mc^2 + \frac{L^2}{m C(r)} \right) \right)$$

$$= x^2 r^2 \left(\frac{1-f}{1+f} \right) \left(\frac{(r_{max}-r)(r-r_{min})}{r_{max}^2} \right) \quad - (10)$$

$$= x^2 r^2 \left(\frac{1-f}{1+f} \right) \left(\frac{(r_{max}-r)(r-r_{min})}{r_{min}^2} \right)$$

The quantities on the right hand side of eq. (10)
are all experimentally measurable.

3) The Carter metric can be interpreted as a perturbation of the Minkowski metric:

$$ds^2 = c^2 dt^2 - dr^2 - r^2 d\theta^2 \quad - (5)$$

and reduces to the Minkowski metric when:

$$AC^{1/2} = BC^{1/2} = 1, \quad C(r) = r^2, \quad - (6)$$

i.e. $A = B = \frac{1}{r}, \quad - (7)$

so eq. (1) becomes:

$$\left(\frac{dr}{d\theta}\right)^2 = r^4 \left(\frac{E^2 - m^2 c^4}{c^2 L^2} - \frac{1}{r^2} \right) \quad - (7)$$

$$= r^4 \left(\left(\frac{p}{L}\right)^2 - \frac{1}{r^2} \right)$$

as in note 233(3).

In this limit:

$$E = \gamma mc^2, \quad p = \gamma mv \quad - (8)$$

$$\text{so } \left(\frac{dr}{d\theta}\right)^2 = r^2 \left(\left(\frac{v}{r\omega}\right)^2 - 1 \right). \quad - (10)$$

For a circular orbit:

$$\frac{dr}{d\theta} = 0. \quad - (11)$$

The above procedure assumes that:

$$r_0 = 0, \quad d_1 = 0. \quad - (12)$$

4) However, the presence of a finite r_0 is important to eliminate unphysical singularities. The original Schwarzschild metric contained a finite r_0 . When this is properly considered there is no big bang and no black holes. So as a first approximation it is reasonable to assume:

$$A = B = \frac{1}{|r-r_0|}, \quad C(r) = |r-r_0|^2 \quad - (13)$$

i.e. $n = 1, \quad d_1 = 0 \quad - (14)$

in eq (6). This means that the whole of the calculation is carried through with r replaced by $|r-r_0|$.

In the Newtonian limit this procedure leads to the result:

$$\left(\frac{v}{|r-r_0| \omega} \right)^2 = 1 + \left(\frac{1+\epsilon}{1-\epsilon} \right) \left(\frac{r_{\max} - |r-r_0|}{r_{\max}^2} (|r-r_0| - r_{\min}) \right) \quad - (15)$$

and that the time processing ellipse is:

$$|r-r_0| = \frac{d}{1+\epsilon \cos(x\theta)} \quad - (16)$$

5) ICB (inter metric):

$$r_0 = \frac{2MG}{c^2} \quad - (17)$$

as given originally by Schwarzschild in his letter to Einstein of 22nd Dec. 1915.

This is a simple illustration of the use of the Schwarzschild/Collins theory. In this case the distance between object m and object M is measured by $|r - r_0|$.

The false Schwarzschild metric is:

$$ds^2 = c^2 dt^2 = \left(1 - \frac{r_0}{r}\right) c^2 dt^2 - \frac{dr^2}{\left(1 - \frac{r_0}{r}\right)} - r^2 d\theta^2 \quad - (18)$$

and there is a singularity at $r = 0$. - (19)

but a simple example of a true Schwarzschild/

inter metric is:

$$ds^2 = c^2 dt^2 = \left(1 - \frac{r_0}{|r - r_0|}\right) c^2 dt^2 - \frac{dr^2}{\left(1 - \frac{r_0}{|r - r_0|}\right)} - \frac{|r - r_0|^2}{|r - r_0|} d\theta^2 \quad - (19)$$

and there is no singularity at $r = 0$.