

233(2): Time precession Ellipse form of Minkowski Metric.

As shown in UFT 149 and UFT 150B the Minkowski metric produces the orbit:

$$\left(\frac{dr}{dt}\right)^2 = r^4 \left(\frac{1}{b^2} - \frac{1}{a^2} - \frac{1}{r^2} \right) \quad (1)$$

where

$$a = \frac{L}{mc}, \quad b = \frac{cL}{E} \quad (2)$$

The so called "Schwarzschild" metric produces the orbit:

$$\left(\frac{dr}{dt}\right)^2 = r^4 \left(\frac{1}{b^2} - \left(1 - \frac{r_0}{r}\right) \left(\frac{1}{a^2} + \frac{1}{r^2} \right) \right) \quad (3)$$

The time precession ellipse is given by:

$$r = \frac{d}{1 + e \cos(x\theta)} \quad (4)$$

where d is the half right latitude, e the eccentricity and x the precession factor. From eq. (4)

$$\frac{dr}{dt} = \frac{d e x \sin(x\theta)}{(1 + e \cos(x\theta))^2} \quad (5)$$

$$= \frac{e x}{d} r^2 \sin \theta$$

$$\text{So } \left(\frac{dr}{dt}\right)^2 = \left(\frac{e x}{d}\right)^2 r^4 \sin^2(x\theta) \quad (6)$$

From eq. (4):

$$2) \quad \cos(x\theta) = \frac{1}{\epsilon} \left(\frac{d}{r} - 1 \right) \quad - (7)$$

Now use:

$$\cos^2(x\theta) + \sin^2(x\theta) = 1 \quad - (8)$$

$$\text{So} \quad \sin^2(x\theta) = 1 - \cos^2(x\theta) \quad - (9)$$

$$= 1 - \frac{1}{\epsilon^2} \left(\frac{d}{r} - 1 \right)^2 \quad - (10)$$

$$\text{and} \quad \left(\frac{dr}{d\theta} \right)^2 = \left(\frac{\epsilon x}{d} \right)^2 r^4 \left(1 - \frac{1}{\epsilon^2} \left(\frac{d}{r} - 1 \right)^2 \right) \quad - (11)$$

From eqs. (1) and (11):

$$x^2 = \frac{\left(\frac{d}{\epsilon} \right)^2 \left(\frac{1}{b^2} - \frac{1}{a^2} - \frac{1}{r^2} \right)}{1 - \frac{1}{\epsilon^2} \left(\frac{d}{r} - 1 \right)^2} \quad - (12)$$

For an ellipse:

$$d = (1 - \epsilon^2) a_0 \quad - (13)$$

$$= (1 + \epsilon) r_{\min}$$

$$= (1 - \epsilon) r_{\max}$$

3) where a_0 is the semi major axis, r_{\min} the distance of closest approach and r_{\max} the distance of maximum separation.

Under these conditions, special relativity produces the true precessing ellipse.

As:

$$r \rightarrow \infty \quad - (14)$$

$$x^2 \rightarrow \frac{d^2}{\epsilon^2 - 1} \left(\frac{1}{b^2} - \frac{1}{a^2} \right) \quad - (15)$$

$$= \frac{d^2}{1 - \epsilon^2} \left(\frac{1}{a^2} - \frac{1}{b^2} \right)$$

$$x^2 \xrightarrow{r \rightarrow \infty} \left(\frac{(1 - \epsilon)^2}{1 - \epsilon^2} \right) r_{\max}^2 \left(\frac{1}{a^2} - \frac{1}{b^2} \right) \quad - (16)$$

For special relativity to reduce to the Newtonian case:

$$x \rightarrow 1 \quad - (17)$$

$$\text{i.e.} \quad \frac{1}{a^2} - \frac{1}{b^2} \xrightarrow{r \rightarrow \infty} \frac{1 - \epsilon^2}{(1 - \epsilon)^2} \frac{1}{r_{\max}^2}$$

$$= \left(\frac{1 - \epsilon}{1 + \epsilon} \right) \frac{1}{r_{\max}^2} \quad - (18)$$