

232(i): The Corret Force Law for a Precessing Ellipse

The force is defined from Lagrangian dynamics

$$\text{as: } \frac{d^2}{dt^2} \left(\frac{1}{r} \right) + \frac{1}{r} = - \frac{mr^2}{L^2} F \quad (1)$$

where m is mass and L is angular momentum. The precessing ellipse is defined by:

$$r = \frac{d}{1 + \epsilon \cos(x\theta)} \quad (2)$$

for small x . As x becomes larger eq. (2) gives parabolic orbits.

By differentiation:

$$\frac{d}{dt} \left(\frac{1}{r} \right) = - \frac{x\epsilon}{d} \sin(x\theta) \quad (3)$$

$$\frac{d^2}{dt^2} \left(\frac{1}{r} \right) = - \frac{x^2\epsilon}{d} \cos(x\theta) \quad (4)$$

Using eqs. (1), (2) and (4):

$$- \frac{x^2}{d} \left(\frac{d}{r} - 1 \right) + \frac{1}{r} = - \frac{mr^2}{L^2} F \quad (5)$$

$$\text{i.e. } F = - \frac{L^2 x^2}{mr^2 d} + \frac{L^2}{mr^3} (x^2 - 1) \quad (6)$$

1) This is not the force law of Einsteinian general relativity. The latter is an incorrect theory:

Using eq. (6) & eq. (1):

$$\begin{aligned}\frac{d^2}{dt^2} \left(\frac{1}{r} \right) + \frac{1}{r} &= -\frac{mr^2}{L^2} \left(\frac{L^2}{mr^3} \frac{d}{dt} \left(\frac{d}{dt} - 1 \right) \frac{-L^2}{mr^3} \right) \\ &= -\frac{d}{dt} \left(\frac{d}{dt} - 1 \right) + \frac{1}{r} \\ &= \frac{d^2}{dt^2} - \frac{d}{dt} + \frac{1}{r} \quad - (7)\end{aligned}$$

So:

$$\boxed{\frac{d^2}{dt^2} \left(\frac{1}{r} \right) + \frac{d}{dt} = \frac{d^2}{dt^2}} \quad - (8)$$

On the other hand, Einsteinian general relativity (EGR) produces:

$$\frac{d^2}{dt^2} \left(\frac{1}{r} \right) + \frac{1}{r} = A + \frac{B}{r^2} \quad - (9)$$

where A and B are constants. Eq. (8) is compatible with eq. (2) but eq. (9) is not.

This is a clear refutation of EGR.

3) If eq. (9) is integrated numerically it will not produce eq. (2). With contemporary computers it is quite simple to integrate eq. (9), but in a textbook such as Maria and Thonka an approximate function is used:

$$u = \frac{1}{r} = A(1 + \epsilon \cos \theta) + BA^2 \epsilon \theta \sin \theta - (10) \\ + BA^2 \left(1 + \frac{\epsilon^2}{2}\right) - \frac{BA^2 \epsilon^2 \cos \theta}{6}$$

This is not the equation of a precessing ellipse, and this can be shown graphically. The precessing ellipse is given by eq. (2) i.e.:

$$u = \frac{1}{r} = \frac{1}{d} (1 + \epsilon \cos(x\theta)) - (11)$$

Maria and Thonka force eq. (10) to become eq. (11). This is a very dubious procedure because the second two terms on the RHS of eq. (10) are just dropped. Eq. (10) is already an approximation to the correct integration of eq. (9). So it is asserted that:

$$u = \frac{1}{r} = ? A(1 + \epsilon \cos \theta) + BA^2 \epsilon \theta \sin \theta - (12)$$

4) Now let us use the further approximation:

$$\begin{aligned} & 1 + \epsilon \cos(\theta - AB\theta) \\ &= 1 + \epsilon (\cos\theta \cos(AB\theta) + \sin\theta \sin(AB\theta)) \\ &\sim 1 + \epsilon \cos\theta + AB\epsilon \sin\theta \quad - (13) \end{aligned}$$

i.e. $\cos(AB\theta) \sim 1$, $\sin(AB\theta) \sim AB\theta$ - (14)

This is true only if $AB\theta \ll 1$. - (15)

In this approximation:

$$\frac{1}{r} \sim \frac{1}{d} (1 + \epsilon \cos((1-AB)\theta)) \quad - (16)$$

Comparing eqs. (2) and (16):

$$\boxed{x = 1 - AB} \quad - (17)$$

The precession is:

$$\Delta\theta = 2\pi(1-x) = 2\pi AB \quad - (18)$$

So it is very clear that EGR produces
the precession "by magic."

5) The correct perihelia precession is:

$$\Delta\theta = 2\pi(1-x) - (19)$$

and x must be measured experimentally, it
cannot be predicted.

Graphical Work

Eq. (10) can be graphed through θ while
range of A and B to show that it is completely
different from eq. (2).

Numerical Work

Eq. (9) can be integrated numerically
for all A and B to show that it does not
give a precessing ellipse.
